

and filtered using a narrow (0.16 nm bandwidth) fiber Bragg grating (FBG) working in reflection. This grating selects only the probe wave amplified by SBS. The Raman pump is a Raman Fiber Laser (RFL) emitting at 1455 nm. The power of this laser can be tuned up to 2.4 W. The RFL beam is divided by a calibrated 50/50 coupler in two beams. Each of them will be coupled to the points X and/or Y represented in Fig. 1, depending on the tested experimental configuration for Raman pumping (co-propagating, counter-propagating or bi-directional propagation with respect to the Brillouin pump pulse).

3. Results

In this section we present experimental results obtained with a standard single mode fiber of 75 km. The width of the Brillouin pump pulses is 20 ns in all cases, thus leading to a 2 meter spatial resolution. In order to achieve this goal, all the configurations have been thoroughly optimised. To correctly determine the Brillouin frequency shift, depletion of the Brillouin pump must be made negligible. Therefore, relatively low Brillouin pump powers and gain values (normally 1-2%) must be preferably chosen. Also, the amount of Raman amplification needed is relatively low, just as much as needed to achieve propagation close to transparency. This optimum value of Raman pumping is around 300 mW on each side (total power around 600 mW in the case of bi-directional amplification) in all studied cases. Additionally, for each configuration we have to search for the optimum Brillouin pump and probe powers so as to avoid pump depletion and maximize the measurement range. The probe power is set high enough to ensure that the signal can be correctly measured in detection. The Brillouin pump is set as high as possible but still trying to avoid depletion. Thus, we show the optimum measurement conditions found with these criteria for all the configurations.

Fig. 2 (a) represents the optimum Brillouin gain traces obtained for a frequency of 10.675 GHz. On average, the contrast of the traces is improved in all the Raman configurations. As we can see, the counter-propagating configuration shows the best improvement in contrast with respect to all the other configurations. The loss in contrast at the input end in the co-propagating and bi-directional configurations is due to the reduction in Brillouin pump power that has been introduced to avoid pump depletion. One remarkable feature of the bi-directional Raman amplification scheme is that the trace obtained is quite flat (the dynamic range in detection does not need to be very large). It is important now to observe the improvements obtained in the determination of the Brillouin shift. Fig. 2 (b) shows the Brillouin shift determined with the gain traces shown before. As we can see, in the conventional BOTDA configuration, the uncertainty in the Brillouin shift grows towards the end of the fiber (up to ± 15 MHz), while it remains bounded in all the Raman configurations. In the case of counter-propagating Raman pump, the deviation between measurements is below ± 3 MHz. This means that the gain enhancement provided by Raman amplification not only results in an increase of the measurement range, but is also particularly useful in reducing the uncertainty of the measurement.

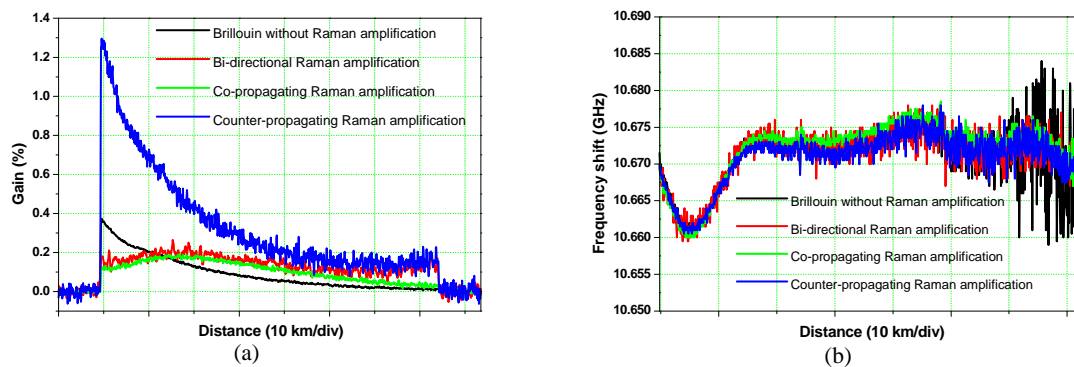


Fig. 2. (a) Brillouin gain for the optimum settings in all the configurations. (b) Brillouin frequency shift of the fiber obtained with the different configurations.

4. References

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