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Interaction of a pulsating vortex rope with the local velocity field in a Francis turbine draft tube

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Abstract. Acoustic resonances in Francis turbines often define undesirable limitations to their operating ranges at high load. The knowledge of the mechanisms governing the onset and the sustenance of these instabilities in the swirling flow leaving the runner is essential for the development of a reliable hydroacoustic model for the prediction of system stability. The present work seeks to study experimentally the unstable draft tube flow by conducting a series of measurements on a reduced Francis Turbine model. The key physical parameters and their interaction with the hydraulic and mechanical system are studied and quantified. In particular, the evolution of the axial and tangential velocity components in the draft tube cone is analysed by means of Laser Doppler Anemometry. Combined with the calculation of the instantaneous vortex rope volume based on flow visualization and the measurement of the pressure fluctuations, the nature of the auto-oscillation in the draft tube flow is investigated.

1. Introduction

Flow instabilities involving cavitation interact dynamically with the surrounding system and may excite acoustic and structural resonances. A problematic example of this surge phenomenon in turbomachines is observed in liquid propellant rockets in the 1960s [1]. This case is at the origin of the development of analytical models identifying the key parameters governing stability by Brennen et al. [2]. Further developments investigating the respective roles of the various model parameters and stability mappings are presented by Tsujimoto et al.[4]. The mass flow gain factor representing the excitation mass source is identified as a key player for stability. Auto-oscillation in diffusers of hydraulic machines and its cause is studied in particular by Chen et al.[7], Alligné et al.[8] and Dörfler et al.[10], choosing different approaches modelling the physics of the draft tube flow. Yonezawa et al. [11] identify the destabilizing effect of the conically shaped diffuser as the main cause for instability.

In a former paper[12] the mass flow gain factor was determined together with the cavitation compliance experimentally in a quasi-static approach by comparing different operating points. To gain better insight into the physics of the full load instability a local approach is chosen in order to study the evolution of various flow and system parameters for an unstable operating point at two different stream wise locations of the conically shaped draft tube cone. The measured quantities include the volume of the vortex rope, the wall pressure in the draft tube cone and the local axial and tangential velocity components.

2. Experimental test case

The tests were performed on a reduced Francis turbine model of specific speed= 0.27 on the EPFL test rig PF3. The top of Figure 1 shows the draft tube cone of the model with the two-component optical probe for the Laser Doppler Anemometry (LDA). The intersection of the laser beams is situated approximately $0.33 \cdot D$ from the symmetry axis of the cone and marks the position of the control volume where the axial and tangential velocity components are measured. Two stream wise measurement positions are chosen and are henceforth referred to as position 1 and position 2 on an upper and lower level of the draft tube cone, respectively. The two pressure sensors for the wall pressure measurements are located on the same horizontal planes as the LDA control volumes, $0.39 \cdot D$ and $1.02 \cdot D$ downstream the runner outlet, respectively. Both pressure sensors share the same angular position in the draft tube cone. The bottom of Figure 1 shows a cut through the upper measurement plane as well as the position of the high speed camera for the flow visualization and the LED backlight.

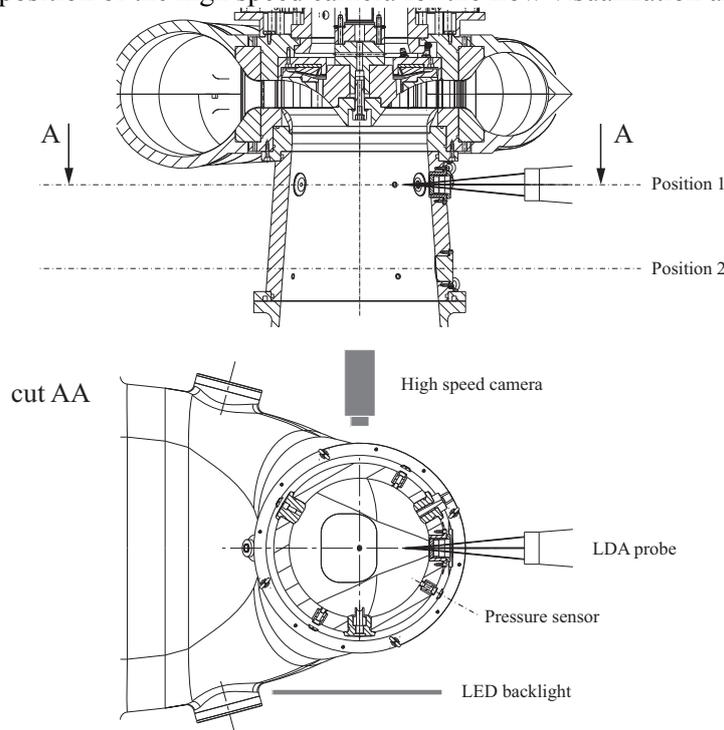


Figure 1. Instrumented draft tube cone with upper measurement plane

3. Measurement parameters

The measurements of the wall pressure, the velocity components and the visualization of the vortex rope are carried out simultaneously by sharing the same trigger between the different measurement systems. Piezo-resistive pressure sensors are used to measure the wall pressure in the draft tube cone. The axial and circumferential velocity components are measured with a state of the art LDA system in non-coincidence mode and the water is seeded with hollow glass sphere particles. A high speed camera is used to film the vortex rope. The uniform LED backlight source provides good contrast between the liquid and gaseous phase. The operating conditions for the investigated instability are summarized in

Table 1.

Table 1. Summary of operating conditions

Position	EM (J kg ⁻¹)	QM (m ³ s ⁻¹)	NED (-)	QED (-)	(-)	Fr (-)
1	115.2	0.37	0.275	0.279	0.135	5.8
2	113.1	0.36	0.277	0.275	0.118	5.8

As it can be seen, the operating conditions for the two separately performed measurements at position 1 and 2, respectively, are subject to small variations. This is due to the fact that the stability conditions may slightly vary between two measurements and so certain operating conditions have to be tuned in order to provoke the instability.

4. Results and analysis

The pulsation of the vortex rope during one cycle of the instability is documented in Figure 2. The six different stages taken from a video recorded during the LDA measurement at position 1 are later identified with the letters A to F. The arrow indicates the location of the LDA measurement. For the calculation of the cavity volume based on its local diameter, the hypothesis of an axis-symmetric vortex rope is made. Unlike in [12], the use of a filter for the edge detection of the vortex rope was not necessary for the major part of the vortex rope due to the good contrast between the liquid and the gaseous phase and simple thresholding was applied. Optical deformations due to the conical shape of the draft tube are not taken into account. The accuracy of the volume's approximation varies furthermore with the rope shape. During the repetitive formation of the cavity, for instance, it is not possible to define clear edges, as can be seen in picture C of Figure 2.

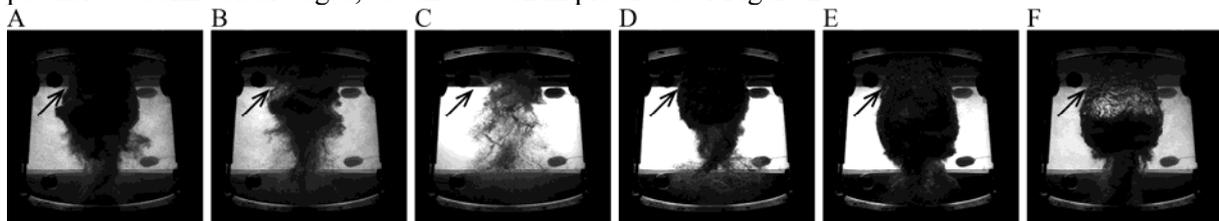


Figure 2. Vortex rope during one cycle of the full load instability

The evolution of the instantaneous vortex rope volume V_C , the wall pressure and the axial velocity C_m in the meridional plane as well as the tangential velocity C_U during four full cycles is shown in Figure 3 for position 1. The void fraction is obtained by dividing the approximate instantaneous vortex rope volume V_C by the total volume of the draft tube cone. The pressure coefficient C_P is calculated by dividing the wall pressure signal by the water density and the overall specific energy E^M at the given operating point. The axial velocity C_m is made non-dimensional using the discharge velocity C_Q of the flow across the section of the cone defined by the horizontal measurement plane. The tangential velocity C_U is made non-dimensional using the circumferential speed U of the runner at its outer downstream diameter.

As expected, the wall pressure rises to its maximum when the vortex rope volume decreases to its minimum. The meridional velocity component C_m decreases when the vortex rope volume V_C decreases and appears to behave symmetrically with respect to the pressure coefficient signal C_P . The tangential velocity component C_U increases when the vortex rope volume drops and the pressure coefficient rises, however with a significant phase shift with respect to the meridional component C_m .

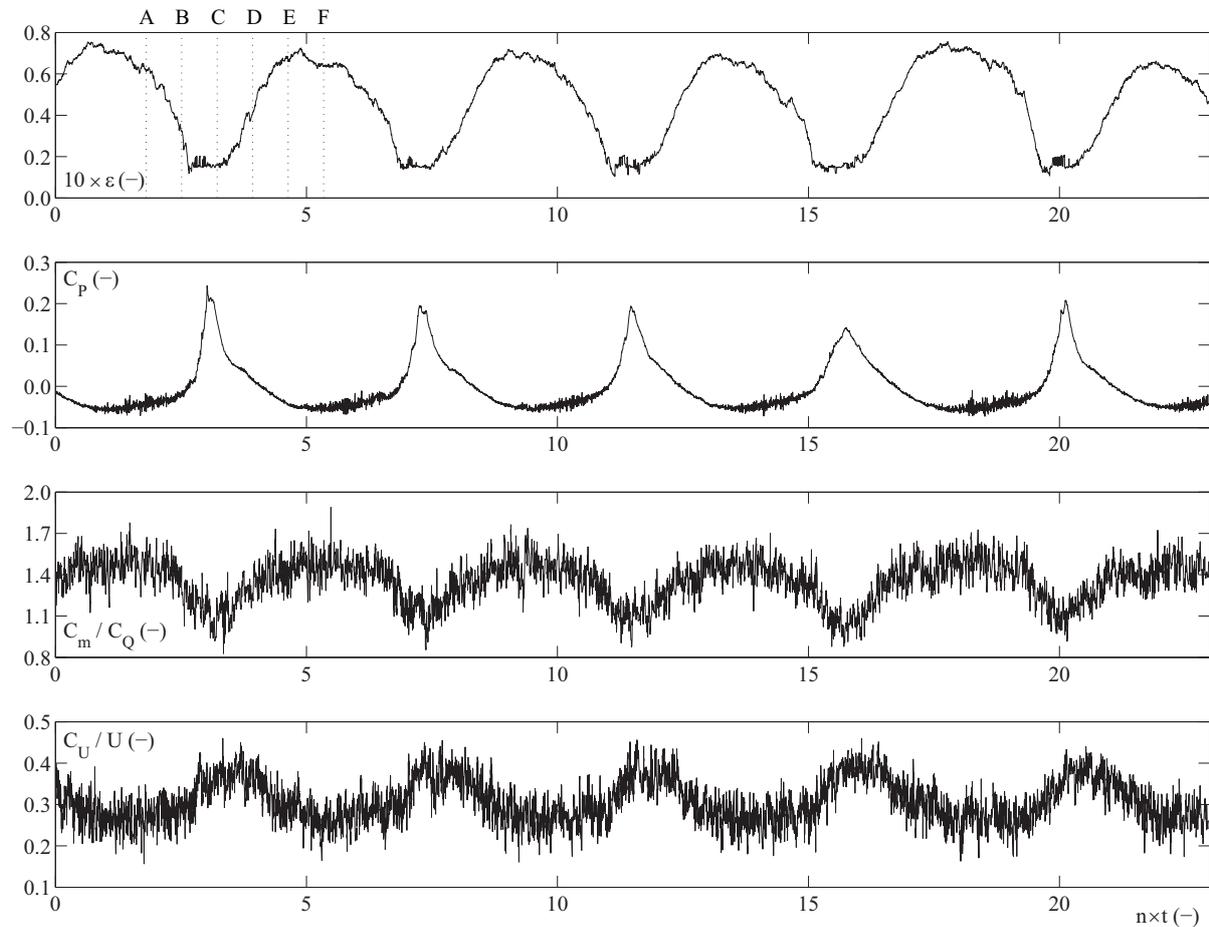


Figure 3. Approximate void fraction and raw signals of the pressure coefficient and the normalized meridional and axial velocity components in the draft tube cone as a function of the number of runner revolutions at position 1

In Figure 4 the C_m and C_U values are averaged with respect to the period of the wall pressure signal for both measurement positions 1, to the left, and 2, to the right. An analytic signal of C_p is therefore calculated, consisting of a real part containing the original data and an imaginary part containing its Hilbert transform [13]. The argument of this analytic signal may be interpreted as an instant phase of the original data and the pressure signal can be divided into its different periods. The velocity signals are then split up individually so that every given fraction of the LDA signal would fit a single period of the C_p signal. The recorded cycles of C_m and C_U are then superposed in one graph together with the superposed C_p cycles. The calculated mean values of C_m and C_U are represented as straight thin lines in Figure 4 and the mean C_p values are represented as straight bold lines, together with the standard deviations. One cycle is divided into 180 sub-cycles of 2 degrees each.

The axial speed C_m at position 1 behaves almost symmetrically with respect to the pressure coefficient C_p and reaches its minimum shortly after the maximum of C_p , as can be seen in the subplot to the top left of Figure 4. Further downstream at position 2 this effect seems to be cancelled out and only a small decrease in C_m is detectable well after the pressure maximum, which is shown in the subplot to the top right. The tangential velocity C_U is significantly elevated at both measurement positions with a phase shift compared to the C_p signal. This is shown in the lower subplots of Figure 4, to the left for position 1 and to the right for position 2.

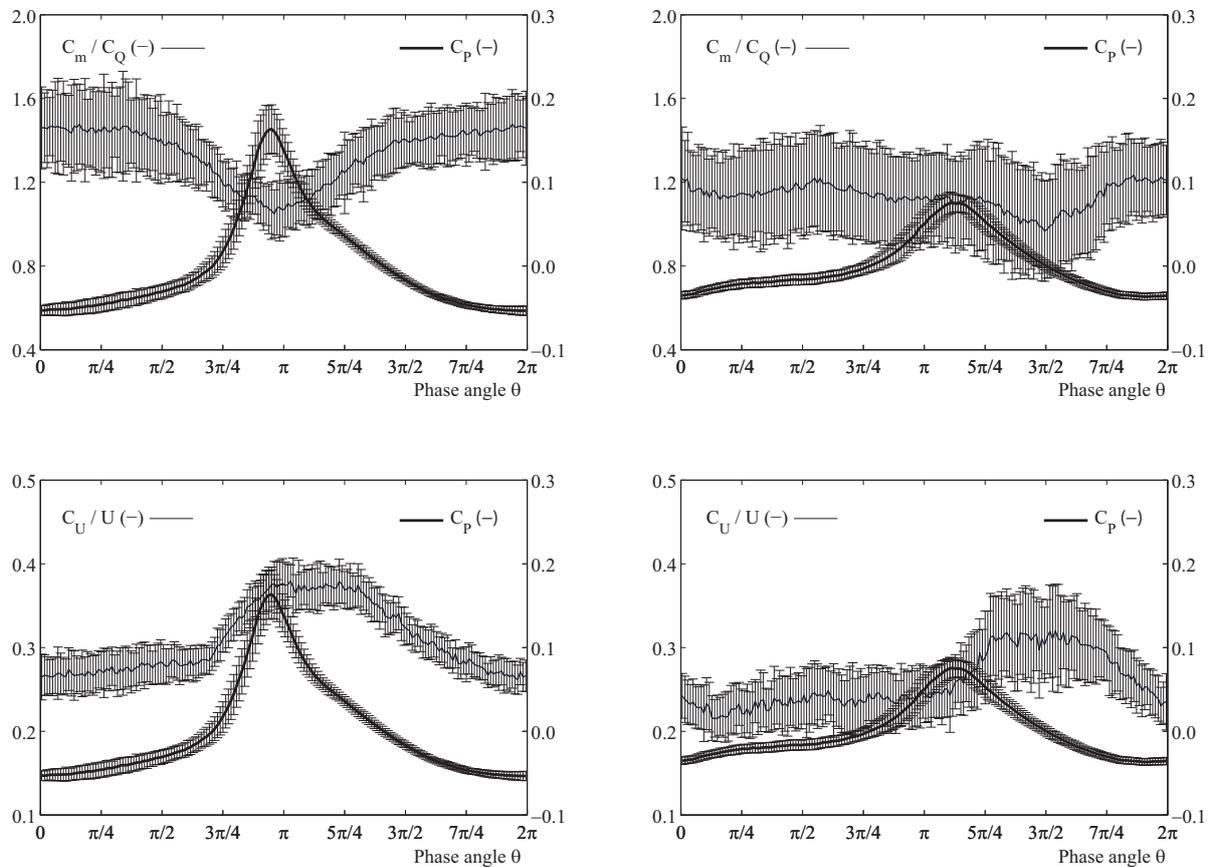


Figure 4. Phase averaged pressure coefficient as well as axial and tangential velocity components with mean and standard deviation values for position 1 to the left and position 2 to the right

5. Conclusions and perspective

The evolution of the axial and tangential velocity components is studied at an unstable full load operating point at two stream wise locations between the inner draft tube cone wall and the pulsating vortex rope. A periodical variation of the axial velocity component C_m is observed at position 1 in form of a significant decrease, whereas this effect is flattened out downstream at position 2. The tangential velocity component C_U periodically increases and decreases at both measurement positions with a delay compared to the C_p signal.

The present analysis is limited to one radial position, $0.33 \cdot D$ from the symmetry axis of the draft tube cone. However, LDA measurements were conducted at different radial positions. The resulting axial velocity components are to be used to calculate the instantaneous discharge at both vertical positions of the draft tube cone and hence the mass flow gain factor of the 1-D model proposed in [8]. Furthermore, a total of twenty-one pressure sensors was distributed throughout the system at the upstream feeding pipe, the spiral case, the draft tube cone, the elbow and the diffuser. The corresponding pressure signals and their dependence on parameters such as the discharge factor Q_{ED} , the head factor n_{ED} and the cavitation number is to be investigated. The instantaneous runner speed and torque were recorded simultaneously with the pressure and LDA signals in order to establish a hydraulic and mechanical energy balance at the instability.

Nomenclature

C_C	Cavitation compliance (m^2)	H	Hydraulic head (m)
C_m	Meridional (axial) velocity ($\text{m}\cdot\text{s}^{-1}$)	n	Runner rotational frequency (s^{-1})
C_P	Pressure coefficient (-)		Specific speed (-)
C_U	Circumferential (tangential) velocity ($\text{m}\cdot\text{s}^{-1}$)	N_{ED}	Speed factor (-)
C_Q	Discharge velocity ($\text{m}\cdot\text{s}^{-1}$)	Q^M	Model volumetric discharge ($\text{m}^3\cdot\text{s}^{-1}$)
D	Runner outlet diameter (m)	Q_{ED}	Discharge factor (-)
E^M	Model specific energy ($\text{J}\cdot\text{kg}^{-1}$)	t	Time (s)
	Void fraction (-)	U	Circumferential runner speed ($\text{m}\cdot\text{s}^{-1}$)
Fr	Froude number (-)	V_C	Cavity volume (m^3)
g	Gravitational acceleration ($\text{m}\cdot\text{s}^{-2}$)		Mass flow gain factor (s)

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